Analytical models for the mooring stiffness of FOWT's and FOWF's

Workshop on dynamics of wind turbines - Universidade de Brasília

Associate Professor Guilherme Rosa Franzini Offshore Mechanics Laboratory Departament of Structural and Geotechnical Engineering Escola Politécnica, University of São Paulo, Brazil

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Objectives

- To present analytical models for evaluating the stiffness associated with the mooring system in floating offshore wind turbines (FOWT's) and floating offshore wind farms (FOWF's);
- To present some results of interest regarding the dynamics of these systems.

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PhD Candidate Giovanni Amaral

- Giovanni has worked on the theme since his last semester as a graduate student (2017);
- Strong collaboration with Prof. Celso Pesce;
- 2018-2020: MSc. dissertation entitled "Analytical assessment of the mooring system stiffness" (see [this link\)](https://doi.org/10.11606/D.3.2020.tde-05112020-114447);
- Results from the MSc. dissertation: Amaral, Pesce & Franzini. Mooring system stiffness: A six-degree-of-freedom closed-form analytical formulation, *Marine Structures*, v. 84, 103189, 2022. The paper can be downloaded from [here;](https://doi.org/10.1016/j.marstruc.2022.103189)
- PhD research: Extension to FOWFs (planar problem) and associated studies.

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Hypotheses

- The floating unit is a rigid body (In some concepts of FOWTs, this may not hold);
- The forces due to the mooring system come from a potential function $V \Rightarrow$ They only depend on the position and aspects such as friction between line and soil are not considered;
- The forces on the line as function of its geometric configuration must be known *a priori*.

Nomenclature

- \bullet $T^{(i)}$: Tension on the i-th mooring line;
- \bullet $F_H^{(i)}$ H ⁽ⁱ⁾, $F_V^{(i)}$ *V* : Vertical and horizontal forces on the i-th mooring line;
- \bullet $h_f^{(i)}$ *f* , *v* (*i*) *f* : anchor-fairlead horizontal and vertical distances;
- \bullet $\hat{e}_h^{(i)}$ $\hat{e}_h^{(i)}$, $\hat{e}_h^{(i)}$ $h_h^{(t)}$: Directional vectors in the vertical and the horizontal directions;
- \bullet $\overrightarrow{P}_{E_n}^{(i)}$ \overrightarrow{P} \overrightarrow{P} \overrightarrow{E}_{ξ} *Eξ* : Position of the i-th fairlead with respect to the fixed and the moving frames;
- \bullet $\overrightarrow{A}^{(i)}_{F_u}$ $E_{\text{r}}^{(t)}$: Anchor position with respect to the fixed frame.

Problem description

Extracted from Amaral *et al*. (2022) - *Mar. Struc.*

The generalized coordinate vector:

$$
\mathbf{q} = [r_x \quad r_y \quad r_z \quad \phi \quad \theta \quad \psi]^t \qquad (1)
$$

where:

$$
\mathbf{r} = [r_x \quad r_y \quad r_z]^t \tag{2}
$$

and

 $\theta = [\phi \quad \theta \quad \psi]^t$ (3)

Geoemetric relations

$$
\overrightarrow{P}_{E_{\xi}}^{(i)} = \left(P^{(i)} - G\right) = \left[p_{\xi}^{(i)} \quad p_{\eta}^{(i)} \quad p_{\zeta}^{(i)}\right]^t \quad (4)
$$

$$
\overrightarrow{P}_{E_x}^{(i)} = \left(P^{(i)} - O\right) = \left[p_x^{(i)} \quad p_y^{(i)} \quad p_z^{(i)}\right]^t = \mathbf{r} + \left[\mathbb{R}\right]_{E_x \mid E_{\xi}} \overrightarrow{P}_{E_{\xi}}^{(i)} \tag{5}
$$

 $E_{\text{X}}^{(i)} = (A^{(i)} - O) = \begin{bmatrix} a_x^{(i)} & a_y^{(i)} & a_z^{(i)} \end{bmatrix}^t$ (6)

$$
h_f^{(i)} = \sqrt{(a_x^{(i)} - p_x^{(i)})^2 + (a_y^{(i)} - p_y^{(i)})^2}
$$
 (7)

$$
n_f^{(i)} = n_f^{(i)} - n_f^{(i)}
$$
 (8)

$$
v_f^{(1)} = p_z^{(1)} - a_z^{(1)} \tag{8}
$$

a (*i*) *^y* − *p* (*i*)

$$
\hat{e}_h^{(i)} = \cos \alpha^{(i)} \hat{e}_x + \sin \alpha^{(i)} \hat{e}_y \tag{9}
$$

$$
\hat{e}_v^{(i)} = -\hat{e}_z \tag{10}
$$

where:

$$
\cos \alpha^{(i)} = \frac{a_x^{(i)} - p_x^{(i)}}{r^{(i)}}
$$
(11)

 $\overrightarrow{A}^{(i)}_{F}$

Forces on the mooring lines

• Tension tension on the i-th mooring line:

$$
\overrightarrow{T}^{(i)}(h_f^{(i)}, v_f^{(i)}) = F_H^{(i)}(h_f^{(i)}, v_f^{(i)})\hat{e}_h^{(i)} + F_V^{(i)}(h_f^{(i)}, v_f^{(i)})\hat{e}_v^{(i)}
$$
(13)

where $\mathrm{a}F_H^{(i)}$ $F_H^{(i)}$ and $F_V^{(i)}$ $V_V^{(1)}$ are the horizontal and vertical forces components;

• Vector of generalized forces:

$$
\mathbf{Q} = \begin{bmatrix} Q_{r_x} & Q_{r_y} & Q_{r_z} & Q_{\phi} & Q_{\theta} & Q_{\psi} \end{bmatrix}^t
$$
(14)

$$
Q_j = \sum_{i=1}^N Q_j^{(i)} = \sum_{i=1}^N \overrightarrow{T^{(i)}} \cdot \frac{\partial \overrightarrow{P}^{(i)}}{\partial q_j} = \sum_{i=1}^N \left(F_H^{(i)} \hat{e}_h^{(i)} + F_V^{(i)} \hat{e}_v^{(i)} \right) \cdot \frac{\partial \overrightarrow{P}^{(i)}}{\partial q_j}
$$
(15)

Generalized forces

 $Q_{r_x}^{(i)} = F_H^{(i)}$

 $Q_{r_y}^{(i)} = F_H^{(i)}$

 $Q_{r_z}^{(i)} = -F_V^{(i)}$ *V*

H cos *α* (*i*)

^{(*i*})</sup> $\sin \alpha$ ^{(*i*})

(16)
\n
$$
Q_{\phi}^{(i)} = F_H^{(i)} \left(\cos \alpha^{(i)} \frac{\partial p_x^{(i)}}{\partial \phi} + \sin \alpha^{(i)} \frac{\partial p_y^{(i)}}{\partial \phi} \right) - F_V^{(i)} \frac{\partial p_z^{(i)}}{\partial \phi} \quad (19)
$$
\n
$$
Q_{\theta}^{(i)} = F_H^{(i)} \left(\cos \alpha^{(i)} \frac{\partial p_x^{(i)}}{\partial \theta} + \sin \alpha^{(i)} \frac{\partial p_y^{(i)}}{\partial \theta} \right) - F_V^{(i)} \frac{\partial p_z^{(i)}}{\partial \theta} \quad (20)
$$
\n
$$
Q_{\psi}^{(i)} = F_H^{(i)} \left(\cos \alpha^{(i)} \frac{\partial p_x^{(i)}}{\partial \phi} + \sin \alpha^{(i)} \frac{\partial p_y^{(i)}}{\partial \phi} \right) \qquad (21)
$$

Using Analytical Mechanics

• The generalized forces: Obtained from the potential function $V = V(\mathbf{q}, \Pi)$, Π are parameters of the mooring system (dimensions, etc):

$$
Q_j = -\frac{\partial V}{\partial q_j} \tag{22}
$$

• Stiffness matrix ⇒ Hessian of the potential energy, evaluated at a position **q** 0 :

$$
\mathbb{K}(\mathbf{q}^0) = \left[\frac{\partial^2 V}{\partial q_j \partial q_k}\right]_{\mathbf{q}^0} = -\left[\frac{\partial Q_j}{\partial q_k}\right]_{\mathbf{q}^0}
$$
(23)

After a lot of Algebraic work

$$
K_{jk} = \sum_{i=1}^{N} K_{jk}^{(i)} = -\sum_{i=1}^{N} \left(\frac{\partial F_H^{(i)}}{\partial h_f^{(i)}} \frac{\partial h_f^{(i)}}{\partial q_k} + \frac{\partial F_H^{(i)}}{\partial v_f^{(i)}} \frac{\partial v_f^{(i)}}{\partial q_k} \right) \hat{e}_h^{(i)} \cdot \frac{\partial \vec{P}^{(i)}}{\partial q_j} + \frac{\partial F_V^{(i)}}{\partial v_f^{(i)}} \frac{\partial v_f^{(i)}}{\partial q_k} + \frac{\partial F_V^{(i)}}{\partial v_f^{(i)}} \frac{\partial v_f^{(i)}}{\partial q_k} \right) \hat{e}_v^{(i)} \cdot \frac{\partial \vec{P}^{(i)}}{\partial q_j} + \frac{\partial F_V^{(i)}}{\partial z_j^{(i)}} \frac{\partial v_f^{(i)}}{\partial q_k} \cdot \frac{\partial \vec{P}^{(i)}}{\partial q_j} + \sum_{i=1}^{N} \left[F_H^{(i)} \frac{\partial}{\partial q_k} \left(\hat{e}_h^{(i)} \cdot \frac{\partial \vec{P}^{(i)}}{\partial q_j} \right) + F_V^{(i)} \frac{\partial}{\partial q_k} \left(\hat{e}_v^{(i)} \cdot \frac{\partial \vec{P}^{(i)}}{\partial q_j} \right) \right]
$$
(24)

with $K_{jk}^{(i)}$ being the stiffness coefficient for the i-th mooring line.

After more algebraic work and considering a perfectly polar symmetry

$$
K_{11} = \frac{n}{2}(k_{HH} + \bar{k}_{HH})
$$
\n(25)
\n
$$
K_{12} = K_{51} = \frac{n}{2}(k_{VH}l + k_{HH}p_{\zeta} + \bar{k}_{HH}p_{\zeta})
$$
\n(26)
\n
$$
K_{22} = \frac{n}{2}(k_{HH} + \bar{k}_{HH})
$$
\n(27)
\n
$$
K_{32} = K_{42} = -\frac{n}{2}(k_{VH}l + k_{HH}p_{\zeta} + \bar{k}_{HH}p_{\zeta})
$$
\n(28)
\n
$$
K_{55} = \frac{n}{2}(p_{\zeta}^2k_{HH} + p_{\zeta}^2\bar{k}_{HH} + 2p_{\zeta}lk_{HV} + l^2k_{VV}
$$
\n(30)
\n
$$
K_{55} = \frac{n}{2}(p_{\zeta}^2k_{HH} + p_{\zeta}^2\bar{k}_{HH} + 2p_{\zeta}lk_{HV} + l^2k_{VV}
$$
\n(31)
\n
$$
K_{66} = n\bar{k}_{HH}l^2\left(1 - \frac{h}{l}\right)
$$
\n(32)
\n(33)
\n(34)
\n(35)
\n
$$
K_{66} = n\bar{k}_{HH}l^2\left(1 - \frac{h}{l}\right)
$$
\n(35)

l: radius from the platform vertical central line to the fairleads

Step-by-step procedure for the mooring system stiffness calculation

Extracted from Amaral *et al*. (2022) - *Mar. Struc.*

OC4 Mooring System - Extracted from Amaral et al. (2022) - Mar. Struc.

OC4 Mooring System

Tension functions

• For an extensible "catenary", F_H and F_V are related to h_f and v_f in closed form:

$$
h_f = l - \frac{1}{\gamma} F_V + \frac{l}{EA} F_H + \frac{F_H}{\gamma} \ln \left(\frac{F_V + \sqrt{F_H^2 + F_V^2}}{F_H} \right)
$$
(33)

$$
v_f = \frac{1}{\gamma} \sqrt{F_H^2 + F_V^2} + \frac{1}{2} \frac{F_V^2}{EA\gamma}
$$
(34)

• Newthon-Raphson scheme is adopted for evaluating the stiffness.

Stiffness matrices at neutral position: Analytical and numerical results

Stiffness matrices

- Numerical results: Very small displacements (????) are prescribed and the forces on the floating unit are computed;
- Numerical results: Coefficients of the stiffness matrices are computed by numerical differentiation;
- Numerical results: High computational cost.
- Analytical results: Practically no computational cost;
- Analytical results: Very useull in the early stages of design, in which a large number of conditions must be considered;
- Analytical results: Valid under the modeling hypotheses.

Influence of the pretension

Extracted from Amaral *et al*. (2022) - *Mar. Struc.*

- In the plot at left, both the stiffness and the pretension are normalized with respect to the nominal values;
- These information are relevant in the design stages and can be easily obtained from the analytical formulation.

Non-neutral yaw position

- Consider that a certain load lead to a different static yaw position $\psi = 5^{\circ}$
- The new stiffness matrix, obtained from the linearization around the static position is

$$
[\mathbb{K}]_{\tilde{\psi}} = \left(\begin{array}{ccccc} 7.16E+1 & 0 & 0 & -9.14E+1 & -1.08E+2 & 0 \\ 0 & 7.16E+1 & 0 & 1.08E+2 & -9.14E+4 & 0 \\ 0 & 0 & 1.92E+1 & 0 & 0 & 1.63E+2 \\ -9.14E+1 & 1.08E+2 & 0 & 8.77E+4 & -5.14E+3 & 0 \\ -1.08E+2 & -9.14E+4 & 0 & -5.14E+3 & 8.77E+4 & 0 \\ 0 & 0 & 1.63E+2 & 0 & 0 & 1.19E+5 \end{array}\right) (37)
$$

Stiffness coefficients at different positions - Amaral et al. (2022) - Mar. Struc.

- Coefficients *K*11, *K*12, *K*13, *K*22, *K*23, *K*³³ as functions of *r* ∗ *^x* and *r* ∗ *y* ;
- *r* ∗ *x*,*y* : Surge/Sway displacements, normalized with respect to the fairlead radius;
- Unities: kN, kNm and rad.

Stiffness coefficients at different positions - Amaral et al. (2022) - Mar. Struc.

- Coefficients *K*14, *K*15, *K*16, *K*24, *K*25, *K*²⁶ as functions of *r* ∗ *^x* and *r* ∗ *y* ;
- *r* ∗ *x*,*y* : Surge/Sway displacements, normalized with respect to the fairlead radius;
- Unities: kN, kNm and rad.

Stiffness coefficients at different positions - Amaral et al. (2022) - Mar. Struc.

- Coefficients *K*44, *K*45, *K*46, *K*55, *K*56, *K*⁶⁶ as functions of *r* ∗ *^x* and *r* ∗ *y* ;
- *r* ∗ *x*,*y* : Surge/Sway displacements, normalized with respect to the fairlead radius;
- Unities: kN, kNm and rad.

Free vibrations - Undamped natural periods and mode shapes

- In the formulation herein, the mooring system accounts only for the stiffness matrix;
- Damping effects \Rightarrow can be included using Rayleigh's model;
- Mass matrix: The added inertia matrix must be added to the "physical" mass matrix;
- Added inertia matrix \Rightarrow depend on the wave period. For the motions on the horizontal plane (more affected by the mooring system), the results are obtained considering $T \to \infty$.

Free vibrations - First three natural periods [s]

- The above results are obtained at the neutral condition;
- The natural periods T_4 , T_5 , T_6 are those from the motions on the vertical plane;
- First-order loads: In the wave periods $(3 < T < 30 \text{ s})$;
- Second-order loads: Observed only in irregular seas;
- Second-order loads: Depend on hydrodynamic coefficients and may excite the periods on the horizontal plane (slow-drift).

Free vibrations - First three natural periods [s]

• The periods can be obtained as functions of the offset caused, for example, by an enviromental load.

Extracted from Amaral *et al*. (2022) - *Mar. Struc.*

Free vibrations - Modes on the horizontal plane

Extracted from Amaral *et al*. (2022) - *Mar. Struc.*

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Contextualization

- When the water depth increases, the cost of the mooring system increases ⇒ Shared mooring systems may be interesting for reducing costs;
- Results obtained by Giovanni as part of his PhD thesis \Rightarrow Extension of the formulation previously presented. Only the planar problem is herein considered;
- The same hypotheses previously adopted are considered;
- Hereafter, all the results are extracted from Giovanni's qualifying written report - Amaral, G.A. *"Analytical tools for design and analysis of shared mooring systems for floating wind farms"*, Escola Politécnica. The exam will take place by the end of the first semester of 2023.

Extracted from Amaral (2023).

Modal analysis at the trivial condition (unloaded system)

Modal analysis at the trivial condition (unloaded system)

Modal analysis at the trivial condition (unloaded system)

Natural periods as functions of the average force and incidence

Allows a "big picture" of scenarios in a single plot.

Possibility of internal resonances

• Internal resonances of type *Tm*/*Tn* being a rational number can be investigated ⇒ Consequences are under study!!!!!

$$
\bar{F} = 520 \text{ kN}, \ \tilde{F} = 100 \text{ kN}, \text{ and } \alpha = 0 \text{ deg}
$$

Forces in the *x* direction are applied to both platforms: $F(t) = \bar{F} + \tilde{F} \sin(2\pi/Tt)$

• Nonlinear time-domain method leads to equilibrium position larger than the linear results \Rightarrow Under investigation, but this point as been obtained by using the MMTS to the Helmholtz-Duffing 1DoF oscillator;

$$
\bar{F} = 520 \text{ kN}, \ \tilde{F} = 100 \text{ kN}, \text{ and } \alpha = 0 \text{ deg}
$$

Forces in the *x* direction are applied to both platforms: : $F(t) = \bar{F} + \tilde{F} \sin(2\pi/Tt)$

- Time-domain method using the stiffness matrix at the trivial condition \Rightarrow Important errors;
- Time-domain method using the "true" stiffness matrix \Rightarrow In agreement with the analytical solution from FRFs;
- Time-domain method using the full nonlinear stiffness forces \Rightarrow Good agreement with the linear results for low-period forcing. Differences observed for $T > 100$ s.

 \bar{F} = 520 kN, \tilde{F} = 100 kN, T = 92 s and α = 0 deg

Forces in the *x* direction are applied to both platform: $F(t) = \bar{F} + \tilde{F} \sin(2\pi/Tt)$

• In this case, the different methodologies do not reveal important differences;

 \bar{F} = 520 kN, \tilde{F} = 100 kN, $T = 108$ s and $\alpha = 0$ deg

Forces in the *x* direction are applied to both platform: $F(t) = \bar{F} + \tilde{F} \sin(2\pi/Tt)$

- Important differences appear;
- The nonlinear (quadratic) hydrodynamic damping the appears in the mooring linear is not incorporated into the quasi-static nonlinear mooring model;

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Final remarks

- Analytical models for the mooring system stiffness were/are beign developed;
- Analytical models: Fully compared with FEM models for the 1 platform case. For FOWFs, the first results indicate very good agreement in terms of natural periods;
- Analytical models: A number of scenarios can be analysed in a fraction of seconds. The same analyses can take hours in FEM codes;
- Analytical models: Very useful in the design and analyses of FOWT/FOWFs ⇒ Implemented in a Matlab® code for these purposes (SuSSA - SubSea Systems Analysis);
- Environmental forcing models: We are trying to improve them for enhancing SuSSA

Final remarks

- Analytical models: Allowed for optimization of the mooring configuration \Rightarrow Giovanni did this for his qualifying exam:
- Most of the works on the literature focus on higher-order hierarchical models for the mooring line;
- Usually, the stiffness matrix and natural periods are calculated only on the trivial position \Rightarrow They can significantly vary with the mean forces (and the corresponding offsets);
- Usually, the design of mooring systems consider static limits \Rightarrow The methodology under development can assess dynamic responses with low computational cost;
- Higher-order hierarchical models: Important for accounting cases not covered by the hypotheses of the analytical models \Rightarrow Run selected scenarios for more detailed analysis.

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